Revisiting slow fast dynamics with piecewise linear differential systems.

A.E. Teruel



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FERNÁNDEZ, DESROCHES, KRUPA, A.E.T. Canard solutions in planar piecewise linear systems with three zones Dyn. Sys. 2016. PROHENS, A.E.T. Canard trajectories in 3D piecewise linear systems DCDS, 2013



Contents

- Canard phenomenon. Supercritical Hopf bifurcation
- PWL Canard explosion
- Subcritical Hopf bifurcation
- PWL Saddle-node canard orbits
- Conclusions

Canard orbit, Canard explosion

- Limit cycle flowing close to a repelling invariant manifold.
- Sudden growing of the amplitude of the canard limit cycles.
- Example: Van der Pol system

$$\begin{array}{l} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(a - x) \end{array} \right\}, \quad f(x) = x \left(\frac{x^2}{3} - 1\right), \quad 0 < \varepsilon \ll 1$$



Canard orbit, Canard explosion

- Limit cycle flowing close to a repelling invariant manifold.
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Change of time $\tau = \varepsilon t$

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(a - x) \end{array} \right\} \qquad \begin{array}{c} \varepsilon x' = y - f(x) \\ y' = a - x \end{array} \right\}$$

Considering the limit problems, $\varepsilon = 0$

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = 0 \end{array} \right\} \quad \begin{array}{c} 0 = y - f(x) \\ y' = a - x \end{array} \right\} \quad \begin{array}{c} y' = f_x x' \\ x' = \frac{a - x}{x^2 - 1} \end{array} \right\}$$



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Canard explosion (geometrical approach) Change of time $\tau = \varepsilon t$

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(\mathbf{a} - x) \end{array} \right\} \qquad \begin{array}{c} \varepsilon x' = y - f(x) \\ y' = \mathbf{a} - x \end{array} \right\}$$

Fenichel Theorem, GSP Theory





Canard explosion (geometrical approach) Change of time $\tau = \varepsilon t$

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Fenichel Theorem, GSP Theory

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = 0 \end{array} \right\} \quad J = \left(\begin{array}{c} -f'(x) & 0 \\ 0 & 0 \end{array} \right)$$



Change of time $\tau = \varepsilon t$

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(a - x) \end{array} \right\} \qquad \begin{array}{c} \varepsilon x' = y - f(x) \\ y' = a - x \end{array} \right\}$$

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a=1.epsilon=0

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Change of time $\tau = \varepsilon t$

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S. Fernández-García, M. Desroches, M. Krupa and A.E. Teruel Canard solutions in planar piecewise linear systems with 3 zones, Dyn.Sys 2015

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(a - x) \end{array} \right\}$$

$$\left\{ \begin{array}{c} x + k + 1 - (k + m)\sqrt{\varepsilon} - ma & x < -1 \\ -kx - (k + m)\sqrt{\varepsilon} - ma & x \in (-1, -\sqrt{\varepsilon}) \\ m(x - a) & |x| < \sqrt{\varepsilon} \\ x - \sqrt{\varepsilon} + m(\sqrt{\varepsilon} - a) & x > \sqrt{\varepsilon} \end{array} \right.$$

$$\begin{array}{c} \bigtriangleup = m^2 - 4\varepsilon < 0 \\ k = 1, |m| < 2\sqrt{\varepsilon} \end{array}$$

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- S. Fernández-García, M. Desroches, M. Krupa and A.E. Teruel Canard solutions in planar piecewise linear systems with 3 zones, Dyn.Sys 2015
- Slow manifolds connection at $a = \tilde{a}(\varepsilon, m)$



- S. Fernández-García, M. Desroches, M. Krupa and A.E. Teruel Canard solutions in planar piecewise linear systems with 3 zones, Dyn.Sys 2015
- Slow manifolds connection at $a = \tilde{a}(\varepsilon, m)$
- ► Clossing equations of canard orbits $\Gamma(h, a, m, \varepsilon)$ through $(-\sqrt{\varepsilon}, h)$

$$\left(\begin{array}{c}F(\tau, a, m, \varepsilon)\\G(\tau, a, m, \varepsilon)\end{array}\right) + \left(\begin{array}{c}\xi_1(\tau, a, m, \varepsilon)\\\xi_2(\tau, a, m, \varepsilon)\end{array}\right) = \left(\begin{array}{c}0\\0\end{array}\right).$$



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PWL Canard Explosion



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Subcritical Hopf bifurcation

Example: FitzHugh-Nagumo

$$\begin{array}{c} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon(\mathbf{a} - x - y) \end{array} \right\}, \quad f(x) = x \left(\frac{x^2}{3} - 1\right), \qquad \qquad J = \left(\begin{array}{cc} 1 - (3a)^{\frac{1}{3}} & 1 \\ -\varepsilon & -\varepsilon \end{array}\right) \\ \bigtriangleup = (3a)^{\frac{1}{3}}\varepsilon > 0 \end{array}$$



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$$\begin{array}{l} \dot{x} = y - f(x) \\ \dot{y} = \varepsilon (a - x - y) \end{array} \right\}, \quad f(x) = x \left(\frac{x^2}{3} - 1\right), \qquad \begin{array}{l} J = \left(\begin{array}{cc} 1 - (3a)^{\frac{1}{3}} & 1 \\ -\varepsilon & -\varepsilon \end{array}\right) \\ \bigtriangleup = (3a)^{\frac{1}{3}}\varepsilon > 0 \end{array}$$



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$$\begin{aligned} \dot{x} &= y - f(x) \\ \dot{y} &= \varepsilon(a - x) \end{aligned} \right\} \\ \left\{ \begin{array}{l} x + k + 1 - (k + m)\sqrt{\varepsilon} - ma & x < -1 \\ -kx - (k + m)\sqrt{\varepsilon} - ma & x < -\sqrt{\varepsilon} \\ m(x - a) & |x| < \sqrt{\varepsilon} \\ x - \sqrt{\varepsilon} + m(\sqrt{\varepsilon} - a) & x > \sqrt{\varepsilon} \\ & \Delta &= m^2 - 4\varepsilon < 0, \quad m = \sqrt{\varepsilon} > 0 \end{aligned} \right.$$



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Th. Slow manifolds connection

$$\begin{split} \tilde{a}(k,\sqrt{\varepsilon}) &= -\frac{e^{\frac{\pi}{\sqrt{3}}} - 1}{e^{\frac{\pi}{\sqrt{3}}} + 1}\sqrt{\varepsilon} - \frac{e^{\frac{\pi}{\sqrt{3}}}}{\left(e^{\frac{\pi}{\sqrt{3}}} + 1\right)^2} \left(\frac{1-k^2}{k^2}\right)\varepsilon^{3/2} + O(\varepsilon^2),\\ \tau_C(k,\sqrt{\varepsilon}) &= \frac{2\pi}{\sqrt{3}}\frac{1}{\sqrt{\varepsilon}} - \frac{1+k}{k} + \frac{1-k^2}{2k^2}\sqrt{\varepsilon} + O(\varepsilon^2). \end{split}$$

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Th. Fixed $\varepsilon \ll 1$, for every $0 < h < h_{\mathbf{p}_L}$ there exists $\hat{a}(k, \sqrt{\varepsilon})$ with the same first terms of the Taylor series expansion as $\tilde{a}(k, \sqrt{\varepsilon})$ such that the system exhibits a canard orbit $\Gamma_s(h, \sqrt{\varepsilon})$ or $\Gamma_b(h, \sqrt{\varepsilon})$.

$$\tilde{a}(k,\sqrt{\varepsilon}) = -\frac{e^{\frac{\pi}{\sqrt{3}}} - 1}{e^{\frac{\pi}{\sqrt{3}}} + 1}\sqrt{\varepsilon} - \frac{e^{\frac{\pi}{\sqrt{3}}}}{\left(e^{\frac{\pi}{\sqrt{3}}} + 1\right)^2} \left(\frac{1-k^2}{k^2}\right)\varepsilon^{3/2} + O(\varepsilon^2),$$



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► Condition of non-hyperbolicity on $\Gamma_s(h, \sqrt{\varepsilon}) \rightarrow e^{t_L \tau_L + t_C \tau_C + t_R \tau_R} = 1.$

$$S(h,\sqrt{\varepsilon}) = \left(1 + \frac{h + (\sqrt{\varepsilon} + \lambda_L^s)(\sqrt{\varepsilon} + a)}{(\lambda_L^q - \lambda_L^s)(\sqrt{\varepsilon} + a)}\right)^{\frac{k}{\lambda_L^s}} \\ \left(1 + \frac{(\sqrt{\varepsilon} + \lambda_R^s)(\sqrt{\varepsilon} - a) - h}{(\lambda_R^q - \lambda_R^s)(\sqrt{\varepsilon} - a)}\right)^{\frac{1}{\lambda_R^s}} - e^{\sqrt{\varepsilon}\tau_c}.$$

- **• Th.** Assume $a = \tilde{a}(h, \sqrt{\varepsilon})$ and $\varepsilon \ll 1$.
 - If k > 1, then $S(h, \sqrt{\varepsilon}) \neq 0$.
 - If k < 1 and $0 < h < h_{p_1}$ is a simple of $S(h, \sqrt{\varepsilon})$, then $\Gamma_s(h, \sqrt{\varepsilon})$ is a non-hyperbolic canard orbit.
- **• Th.** Assume $a = \tilde{a}(h, \sqrt{\varepsilon}), \varepsilon \ll 1$, and k < 1.
 - a) $\lim_{h\searrow 0} S(0,\sqrt{\varepsilon}) > 0$ and $\lim_{h\nearrow +\infty} S(h,\sqrt{\varepsilon}) = -e^{\frac{2\pi}{\sqrt{3}}}$.
 - b) Let $h^* > 0$ be a solution of $S(h, \sqrt{\varepsilon}) = 0$, then $\frac{\partial S}{\partial h}\Big|_{(h^*, \sqrt{\varepsilon})} < 0$.
 - c) Let $h^*(k, \sqrt{\varepsilon})$ be the positive solution of $S(h, \sqrt{\varepsilon}) = 0$, then



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• Chosen a canard orbit $\Gamma_b(h, \sqrt{\varepsilon})$.

$$\begin{split} \tilde{\tau}_L &= \frac{1}{\lambda_L^s} \ln \left(1 + \ldots \right) \quad \tau_{LL} = -\frac{1}{\lambda_{LL}^s} \ln \left(1 + \ldots \right), \\ \tau_{RR} &= -\frac{1}{\lambda_R^s} \ln \left(1 + \ldots \right). \end{split}$$

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• **Th.** Assume
$$a = \tilde{a}(h, \sqrt{\varepsilon})$$
 and $\varepsilon \ll 1$.

- If k < 1, then $S(h, \sqrt{\varepsilon}) \neq 0$.
- If k ≥ 1, exists 0 < h_M < h_{PL} s. t. if 0 < h < h_M is a zero simple of S(h, √ε), then Γ_b(h, √ε) is a non-hyperbolic canard orbit.
- ▶ **Th.** Assume $a = \tilde{a}(h, \sqrt{\varepsilon})$, $\varepsilon \ll 1$, and $1 < k < (1 + e^{\pi/2})$.
 - a) $S(0,\sqrt{\varepsilon}) < 0$ and $S(k,\sqrt{\varepsilon}) > 0$.
 - b) Let $0 < h^* < k$ be a solution of $S(h, \sqrt{\varepsilon}) = 0$, then $\frac{\partial S}{\partial h}\Big|_{(h^*, \sqrt{\varepsilon})} > 0$.
 - c) Let $h^*(k,\sqrt{\varepsilon})$ be the positive solution of $S(h,\sqrt{\varepsilon}) = 0$, then





Figure: Amplitude of the saddle-node canard orbit versus parameter k for $\varepsilon = 1e - 6$. The straight line corresponds with the ordinate of the tangent point \mathbf{p}_L , and coincides with the maximum size of a 3-zones saddle-node canard orbit. Therefore, the discolored part of the curve does not correspond with saddle-node canard orbits.



Figure: Amplitude of the saddle-node canard orbit versus parameter k for $\varepsilon = 1e - 6$. The straight line corresponds with the ordinate of the tangent point \mathbf{p}_L , and coincides with the maximum size of a 3-zones saddle-node canard orbit. Therefore, the discolored part of the curve does not correspond with saddle-node canard orbits.

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